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# Estimation of Ionic Impurities in Poly(propylene Glycol) Diacrylate Monomers/Liquid Crystal E7 Mixtures Using Dielectric Spectroscopy

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**Abstract:** The study investigated the effect of the molecular weight of three difunctional poly(propylene glycol) diacrylates on the temperature-dependent ionic conductivity of these monomers and their blends with an eutectic nematic liquid crystal mixture (E7). The results revealed two distinct regions. At low temperatures, ionic conduction can be described by the Vogel–Tamman–Fulcher (VTF) equation, while at high temperatures, the conductivity data follow the prediction of the Arrhenius model. The Arrhenius and VTF parameters and their corresponding activation energies were determined using the least squares method. In addition, a conductivity analysis based on an ionic hopping model is proposed. Estimates of ion concentrations and diffusion constants were calculated. It was found that both the ionic concentration and the diffusion constant decrease with the increase in the molecular weight of the monomers. The static dielectric permittivity decreases in the following order: TPGDA, PPGDA540, and PPGDA900. This can be explained by the higher dipole moment of TPGDA, which is caused by an enhanced volume density of carbonyl groups.

Keywords: poly(propylene glycol) diacrylate; liquid crystal; dielectric spectroscopy; ionic conductivity



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## 1. Introduction

Polymer-dispersed liquid crystal (PDLC) films are composite materials composed of micron-sized LC domains phase-separated by a polymer matrix [1–4]. These materials are still undergoing intensive studies due to their intriguing properties and potential applications, such as optical shutters, intelligent windows, information displays, and holographic devices [1,2,5,6]. Several techniques with which to produce these materials are available, including radiation-induced polymerization and crosslinking. This process begins with the irradiation of a homogeneous mixture of low-molecular-weight LC and reactive monomers. [4,7,8] As the polymer network grows, the thermodynamic miscibility between the network and the LC decreases, resulting in phase separation and the formation of randomly distributed LC domains within the polymer matrix.

Inorganic impurities (ions) present in the initial monomer/LC mixtures can have adverse effects on the resulting PDLC materials. These effects include increased electrical conductivity, which can lead to a loss of electro-optical or optical properties such as long response time, persistence, and screen flickering [9,10]. Therefore, it is necessary to

Crystals **2024**, 14, 286 2 of 15

investigate the dielectric permittivity and electrical conductivity of the initial monomer/LC materials.

Low-frequency dielectric properties have garnered significant attention due to the process of charge carrier transport. Several authors [11–14] have studied ion transport diffusion in liquids using dielectric theory. Raymond et al. [15] demonstrated that the properties of organosiloxane liquids and polymers are highly dependent on conductivity in relation to humidity. Several studies have investigated the ionic conduction of polymer electrolytes [16–18] to develop multifunctional energy-producing composites, such as structural batteries. In particular, researchers have utilized poly(ethylene glycol) (PEG)-based materials [19–21] and have demonstrated their high conductivity. Several authors [22–24] have studied ionic impurities in organic materials, particularly in LCs, using Iwamoto's hopping model [25] to describe their conductivity behavior.

The purpose of this study is to investigate the dielectric permittivity and electrical conductivity of a model system composed of difunctional acrylic monomers and blends of these monomers with a eutectic nematic LC mixture composed of cyanobiphenyl derivatives (commercially known as E7). Three acrylic polypropylene glycol diacrylate (PPGDA) monomers were used, with reactive chains of different molecular weights between their two acrylic end groups. These materials were chosen because they have already been studied in the UMET laboratory and present interesting electro-optical, thermophysical, and mechanical properties [26–28].

The study will investigate the effect of monomer chain length (i.e., molecular weight) and the presence of ionic impurities on static permittivity and ionic conductivity over a wide range of frequencies and temperatures. The electrical conductivity of the samples, which are in the form of films sandwiched between two electrodes, will be analyzed using the Iwamoto model. This model allows for the estimation of ion concentration and diffusion constants. The conductivity data obtained from the investigated monomer and monomer/E7 samples will be analyzed as a function of temperature using the Arrhenius and Vogel–Tamman–Fulcher (VTF) models [29–33].

## 2. Materials and Methods

#### 2.1. Materials

This study utilized three acrylic difunctional propylene glycol-based monomers with identical chemical structures, differing only in their chain lengths, as determined by their molecular weights: tripropylene glycol diacrylate (TPGDA) with  $M_n = 300~g/.mol$ , polypropylene glycol diacrylate with  $M_n = 540~g/mol$  (PPGDA540), and polypropylene glycol diacrylate with  $M_n = 900~g/mol$  (PPGDA900);  $M_n$  represents the number average of the molecular weight, as given by the sales companies (Cray Valley and Sigma-Aldrich, see below).

Samples were prepared by mixing *X* weight–percent (wt.-%) LC with (100–*X*) wt.-% of the monomer. All products were used as received.

The TPGDA material was purchased from Cray Valley Company (Paris, France), and the other monomers (PPGDA540 and PPGDA900) were purchased from Aldrich Chemical Company (Saint-Quentin-Fallavier, France), while the eutectic nematic LC mixture E7 was obtained from Merck Company (Tokyo, Japan).

#### 2.2. Methods

## 2.2.1. Nuclear Magnetic Resonance Spectroscopy

To determine the chemical composition of the monomers and the number of repetition units, we utilized the proton nuclear magnetic resonance technique ( $^{1}$ H-NMR). The materials were dissolved in deuterated solvent, specifically deuterated chloroform (CDCl<sub>3</sub>) ( $^{10}$  mg/mL). The  $^{1}$ H-NMR spectra were recorded on a Bruker AC300 spectrometer (Bruker, Billerica, MA, USA) operating at 300 MHz. The data are reported in ppm relative to tetramethylsilane (TMS).

Crystals **2024**, 14, 286 3 of 15

### 2.2.2. Gel Permeation Chromatography

The monomers were dissolved in THF at a concentration of 3 mg/mL. The gel permeation chromatography (GPC) technique was used to measure the molar masses at  $T = 25\,^{\circ}$ C. Waters apparatus, including a Waters 515 pump, a Waters 717 plus auto-sampler, a differential refractometer Waters 410, and Stryragel columns HR3 and HR 3E (Waters TM, Milford, MA, USA), was used for the measurement. The experiment lasted nearly 1 h with a flow rate of 1 mL/min.

The GPC measurements yield distinct peaks based on the retention time of each component. The molar mass of each monomer sample was determined by utilizing the calibration curve polynomial.

#### 2.2.3. Thermogravimetric Analysis

Thermogravimetric analysis (TGA) was performed using a Perkin Elmer Pyris 1 analyzer (Perkin-Elmer Company, Shelton, CT, USA) with a mass resolution of 1  $\mu g$  on HT platinum plates. The samples, with an average weight of 8 mg, were analyzed under a nitrogen atmosphere with a flow rate of 20 mL/min. A heating ramp of 10  $^{\circ} C/min$  was applied to the samples in the temperature range of 25 to 600  $^{\circ} C$ .

This analysis is informative because it provides data on the critical temperature ( $T_C$ ) of the material at the beginning of the thermal degradation processes.  $T_C$  is considered as the upper temperature limit for DSC analysis.

## 2.2.4. Differential Scanning Calorimetry

The thermal properties of monomer and monomer/LC materials were determined using the differential scanning calorimetry (DSC) technique (see, for example, in [34]). DSC measurements were performed on a Seiko DSC 220C (Seiko Instruments Inc., Chiba, Japan) equipped with a liquid nitrogen system for cooling experiments. The DSC cell was purged with 50 mL/min of nitrogen. Heating and cooling rates of 10 °C/min and 30 °C/min, respectively, were used in the temperature range from -120 °C to +100 °C. The program begins by cooling the sample, followed by three cycles of heating and cooling to account for any thermal events related to the sample's preparation history. The results presented in this work were obtained from the first heating ramps. At least three duplicate samples, with the same composition and prepared independently, were used in each case to verify the reproducibility of the results. The glass transition temperatures ( $T_g$ ) were determined by taking the midpoint of the transition range in the thermograms.

#### 2.2.5. Dielectric Spectroscopy

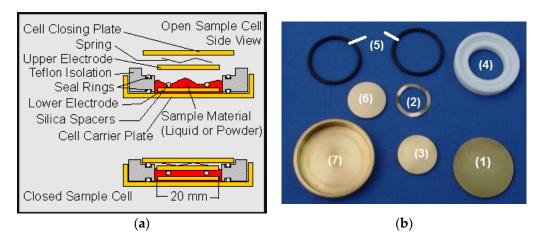
The dielectric measurement spectroscopy was performed using the Concept 80 dielectric spectrometer from Novocontrol  $^{\circledR}$  GmbH (Novocontrol Technologies Company, Montabaur, Germany) (more details are given in [35]). The measurements were performed in the frequency range of 0.1 Hz–1 MHz. The real and imaginary parts of the complex dielectric constant ( $\epsilon^*=\epsilon'-j\epsilon''$ ) were calculated from the measured capacitance  $C_m$  and conductance  $G_m$  at each frequency f:

$$\varepsilon'(f) = \frac{C_{\rm m}(f)}{C_0} \tag{1}$$

$$\varepsilon''(f) = \frac{G_m(f)}{2\pi f C_0} \tag{2}$$

The capacitance of the empty cell  $C_0$  was measured prior to cell filling. The samples were prepared as 50  $\mu$ m thick films and placed between two polished brass electrodes covered with a thin layer of gold (20 mm in diameter) (see Figure 1). They were then inserted into a temperature-controlled sample cell designed for liquid samples. Temperature was regulated using a nitrogen gas cryostat with a stability of better than 0.1 °C.

Crystals **2024**, 14, 286 4 of 15



**Figure 1.** (a) Synoptic diagram of dielectric measurement cell. (b) Photo of the open cell: (1) cell closing plate, (2) spring, (3) upper electrode, (4) Teflon isolation, (5) seal rings, (6) lower electrode, (7) cell carrier plate (reproduced from [36]).

#### 3. Results and Discussion

#### 3.1. Chemical and Structural Properties of Materials

The three materials studied belong to the PPGDA monomer family. The repeating unit consists of a poly(propylene) glycol group. The monomer tripropyleneglycol diacrylate (TPGDA) corresponds to n=3, where n represents the number of repeating units. The other two materials are commercially known as PPGDA540 and PPGDA900.

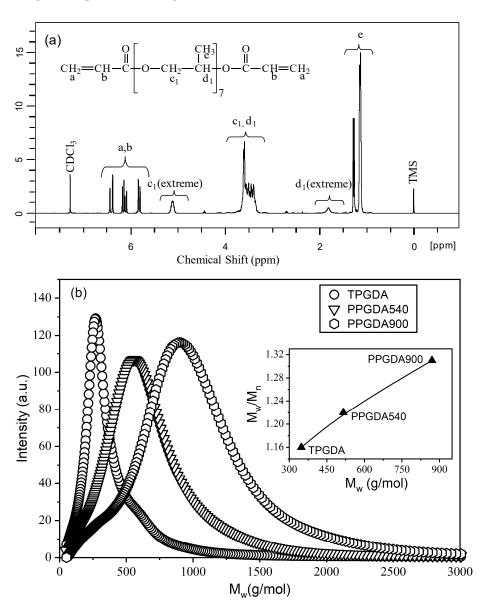
Figure 2a presents the  $^1H$  NMR spectrum obtained from PPGDA540. The spectrum shows three peaks between 5.5 and 6.5 ppm, corresponding to the six acrylic protons (a and b). Protons (d<sub>1</sub> extreme) and (c<sub>1</sub> extreme) appear between 1.5 and 2 ppm and between 4.5 and 5.5 ppm, respectively. The other protons (c<sub>1</sub> and d<sub>1</sub>) and the methyl protons (e) are observed between 3 and 4 ppm and between 1 and 1.5 ppm, respectively. The  $^1H$ -NMR spectra of the TPGDA and PPGDA900 monomers exhibit identical peaks to those shown in Figure 2a. The spectral analysis shows that the TPGDA, PPGDA540, and PPGDA900 monomers correspond to n = 3, 7, and 12, respectively, as determined by calculating and comparing integration ratios.

The molecular weights and molecular weight distributions were determined using the GPC method, in addition to the results obtained by  $^1$ H-NMR spectroscopy. This technique can analyze materials of even lower molecular weight by using the appropriate column material [37]. Figure 2b displays the chromatograms of the three monomers. The measured values  $M_{\rm w}$  of the average molecular weight of the three monomers are 460, 765, and 1120 daltons. Figure 2b shows that the monomers have a polydisperse character in terms of their molecular weight distribution.

Figure 3a presents the TGA analysis results for TPGDA, PPGDA540, and PPGDA900. The figure shows that the TPGDA monomer degrades faster than the other monomers. It remains stable up to 140 °C, after which it degrades to 90% of its initial weight between 195 °C and 325 °C, followed by a second degradation. At 465 °C, the TPGDA monomer is almost completely degraded (2% of the initial mass). The initial degradation may be attributed to lower molar masses (see Figure 2b). The residual material (2%) observed at 465 °C can be linked to the presence of inorganic impurities. It is worth noting that the PPGDA540 monomer degrades at a slower rate than the TPGDA monomer. The substance remains stable up to a temperature of 152 °C, after which it undergoes a weight decrease until it loses 12% of its original weight at 313 °C. The material experiences rapid degradation, with complete degradation occurring at 327 °C. The degradation of 12% between 152 to 313 °C suggests the presence of inorganic impurities in the monomer, although less than in TPGDA monomer. The PPGDA900 monomer is the most stable of the three monomers, undergoing a slow decrease rather than a sudden drop. Its degradation

Crystals **2024**, 14, 286 5 of 15

begins at 221 °C and reaches total degradation at 443 °C. This monomer also contains inorganic impurities (see Figure 3a).



**Figure 2.** (a)  $^{1}$ H NMR spectrum of the monomer PPGDA540 (n = 7), (b) GPC chromatograms of the monomers.

TGA limits the working temperature range for other analyses, such as DSC and dielectric measurements. During this study, the maximum applied temperature did not exceed 100  $^{\circ}$ C. Upon closer examination in Figure 3a (450  $^{\circ}$ C–600  $^{\circ}$ C), it is observed that the concentration of inorganic impurities varies inversely with the molecular weight of the monomers.

Figure 3b presents DSC thermograms for the three monomers, which clearly show glass transitions with different  $T_g$  values. The glass transition temperature of TPGDA, PPGDA540, and PPGDA900 were found to be  $-85\,^{\circ}\text{C}$ ,  $-74\,^{\circ}\text{C}$  and  $-70\,^{\circ}\text{C}$ , respectively. The increase in  $T_g$  with the monomeric chain length is in good agreement with the classical behavior of low-molecular-weight polymeric chains. The relationship between  $T_g$  and  $M_W$  can be described by the well-known Fox–Flory equation [38].

Crystals **2024**, 14, 286 6 of 15

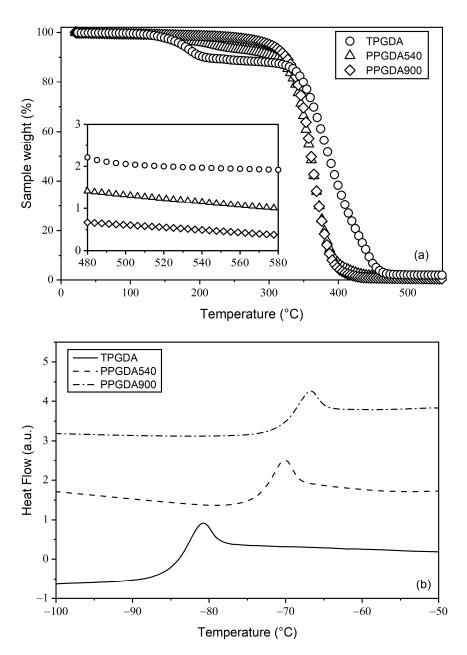


Figure 3. (a) TGA thermograms of the monomers and (b) DSC thermograms of the monomers.

#### 3.2. Dielectric Characterization

As an example, Figure 4a,b present the results of dielectric measurements at room temperature. Two contributions to the dielectric response [39–42] can be distinguished: the first concerns the response of the dipolar moment ( $\epsilon_d^*$ ) and the second is linked to the contribution of ionic impurities (conductivity) ( $\epsilon_{ion}^*$ ). On the ln ( $\epsilon'(f)$ ) plot for TPGDA for example (Figure 4a), two distinct behaviors are evident: at high frequencies (f > 1 kHz),  $\epsilon'$  remains constant and measures the static dielectric constant of the monomer (dipolar moment response), with a value of 7.4. At low frequencies,  $\epsilon'$  increases as f decreases due to the accumulation of ionic charges near the electrodes. The contribution of ions can also be observed on the  $\epsilon''(f)$  plot, specifically in the linear range between 1 Hz and 10 kHz. It is important to note that none of the three monomers exhibit a relaxation process in the studied frequency range at room temperature. This includes the absence of the  $\alpha$  relaxation process due to the low  $T_g$  of these monomers in comparison to room temperature.

Crystals **2024**, 14, 286 7 of 15

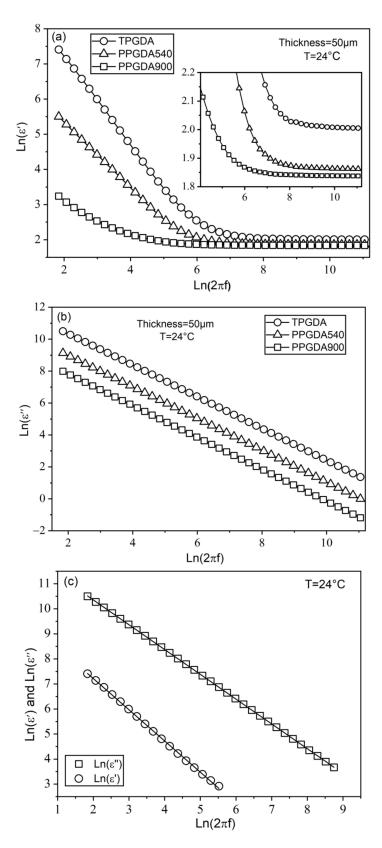
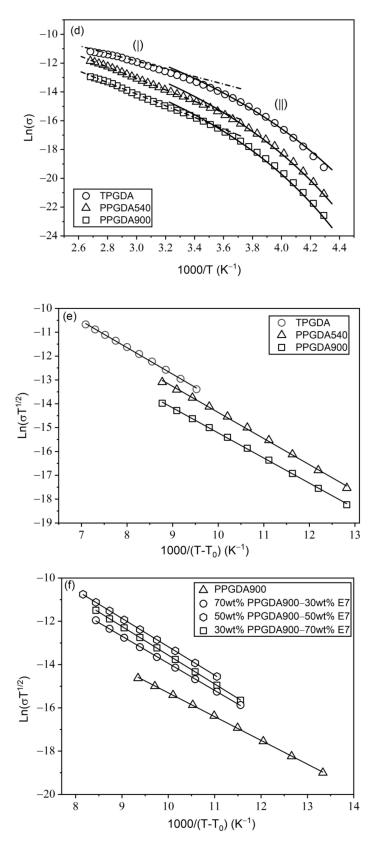


Figure 4. Cont.

Crystals **2024**, 14, 286 8 of 15



**Figure 4.** (a) Low-frequency evolution of  $Ln(\epsilon'(f))$ . (b) Low-frequency evolution of  $Ln(\epsilon'(f))$ , (c) Evolution of permittivities  $\epsilon'$  and  $\epsilon''$  of TPGDA as a function of frequency:  $T=24~^{\circ}C$ , 1~Hz < f < 1~kHz,  $A=2.487\times 10^{-7}$ ,  $B=6.067\times 10^{-7}$ ,  $m=2.319\times 10^{-1}$ ,  $\sigma=2.247\times 10^{-6}~S/m$ . (d) Temperature evolution of the conductivity of the monomers. (e) Conductivity of the three monomers. (f) Conductivity of PPGDA900 monomer/E7 mixtures.

Crystals **2024**, 14, 286 9 of 15

#### 3.3. Conductivity

The dielectric spectra typically describe the ionic contribution using the following formulas [39].

$$\varepsilon'_{\text{ion}} = \frac{A}{(2\pi f)^{m+1}} \left(\frac{\sigma}{\varepsilon_0}\right)^2 \tag{3}$$

$$\varepsilon''_{\text{ion}} = \frac{\sigma}{2\pi f \varepsilon_0} - \frac{B}{(2\pi f)^{m+1}} \left(\frac{\sigma}{\varepsilon_0}\right)^2$$
 (4)

where  $\sigma$  represents the ionic conductivity, and A and B are constants that depend on the interface between monomer and electrodes. They take into account the effects of electrode polarization,  $\epsilon_0$  is the dielectric permittivity of vacuum, and m is a real number between 0 and 1. Figure 4c shows the experimental plots of  $\ln(\epsilon')$  and  $\ln(\epsilon'')$  as a function of  $\ln(2\pi f)$ . When the plot of  $\ln(\epsilon'')$  versus  $\ln(2\pi f)$  yields a linear relationship, and its slope is equal to -1, the conductivity  $\sigma$  will only be ionic (see for example Figures 3.1a and 3.16a in reference [30]). The data in Figure 4b and the top curve in Figure 4c, as representative examples, exhibit slopes equal to -1, which is also the case for all other samples examined. This indicates the absence of electrode polarization effects and other phenomena. As a result, only the contribution of the ionic conductivity in  $\sigma$  is present. Equations (3) and (4) were simultaneously applied in the fitting procedure of the experimental data, allowing for the extraction of the different parameters A, B, m, and  $\sigma$ . On the right-hand side of Equation (4), it was found that the weight of the right term  $((\frac{B}{(2\pi f)^{m+1}} \left(\frac{\sigma}{\epsilon_0}\right)^2)$  is much smaller than that of the left term  $(\frac{\sigma}{2\pi f\epsilon_0})$ .

As a result, Figure 4d presents the Arrhenius plot of the monomers, displaying the temperature-dependent evolution of the calculated ionic conductivity, which increases with temperature. The curves of the three monomers exhibit a curve-like behavior below a certain temperature, which becomes nearly linear above that temperature. The non-linearity of the Arrhenius plot has been reported in several papers and indicates that the ionic transport is correlated with the segmental motion of the polymer chain [40–42]. In this case, the results can be more effectively represented by using the VTF equation (Equation (5)).

$$\sigma = \sigma_1 T^{-\frac{1}{2}} \exp\left(\frac{-E_v}{k(T - T_0)}\right)$$
 (5)

where T is the absolute temperature;  $\sigma_1$ ,  $E_v$  and  $T_0$  are the fitting constants; k represents the Boltzmann constant; and  $\sigma_1$  stands for a pre-exponential factor, which is related to the number of charge carriers.  $E_v$  is the pseudo-activation energy, and  $T_0$  is the critical temperature (ideal glass transition) at which the conductivity becomes zero [43]. Typically,  $T_0$  is set to be 30-50 K below  $T_g$ . The  $T_g$  values were obtained from DSC measurements.  $T_0$  was calculated as  $T_0=T_g-40$  K. Figure 4e,f show the plot of  $\text{Ln}\left(\sigma T^{\frac{1}{2}}\right)$  versus  $\frac{1000}{T-T_0}$  for the three monomers and the PPGDA900/E7 mixtures as examples. The data indicate a strong correlation between the VTF theory and the low-temperature data. However, in the high-temperature region (I), the curve deviates from the VTF equation. The transition from the VTF equation to the Arrhenius equation is clearly visible in this region. This behavior is commonly observed in polymer electrolyte systems. The data in region (I) can be fitted using the Arrhenius equation:

$$\sigma = \sigma_0 \exp\left(\frac{-E_a}{kT}\right) \tag{6}$$

where  $\sigma_0$  is the pre-exponential factor, and  $E_a$  stands for the activation energy of the thermally activated process. The values obtained for the TPGDA monomer are  $\sigma_0 = 0.21 \times 10^{-1} \text{S/m}$  and  $E_a = 0.230$  eV. The fitting parameters obtained from the Arrhenius and VTF equations for the three monomers are shown in Table 1, and those for the monomer/E7 mixtures are presented in Table 2.

Crystals **2024**, 14, 286 10 of 15

Monomer	$\sigma_0(S/m)(\times 10^{-1})$	E <sub>a</sub> (eV)	$\sigma_1\left(S.K^{\frac{1}{2}}.m^{-1}\right)(\times 10^{-2})$	E <sub>v</sub> (eV)	T <sub>0</sub> (K)
TPGDA	0.21	0.230	6.46	0.095	148
PPGDA540	3.15	0.342	3.27	0.094	159
PPGDA900	1.30	0.348	0.91	0.090	163

**Table 1.** Arrhenius and VTF equation fitting parameters of the monomers.

Table 2. Arrhenius and VTF equation fitting parameters of the monomer/E7 mixtures.

Monomer/E7 Mixture	$\sigma_0(S/m)(\times 10^{-1})$	E <sub>a</sub> (eV)	$\sigma_1\bigg(S.K^{\frac{1}{2}}.m^{-1}\bigg)\big(\times 10^{-1}\big)$	E <sub>v</sub> (eV)	T <sub>0</sub> (K)
70wt%TPGDA-30wt%E7	3.33	0.291	5.10	0.104	157
50wt%TPGDA-50wt%E7	11.7	0.304	2.40	0.098	157
30wt%TPGDA-70wt%E7	13.3	0.311	11.70	0.109	157
70wt%PPGDA540-30wt%E7	5.90	0.319	2.24	0.101	162
50wt%PPGDA540-50wt%E7	2.25	0.291	1.15	0.094	163
30wt%PPGDA540-70wt%E7	2.62	0.262	8.95	0.113	163
70wt%PPGDA900-30wt%E7	15.6	0.355	2.22	0.107	162
50wt%PPGDA900-50wt%E7	31.9	0.350	17.4	0.118	166
30wt%PPGDA900-70wt%E7	21.1	0.334	9.60	0.116	165

To complete this section on conductivity, it is assumed that the ionic diffusion phenomena can be described by Iwamoto's non-linear model for ionic hopping [25]. Iwamoto's model considers the presence of a diffusion process and an electrostatic interaction between ions, while also assuming the absence of a space-accumulating charge field at the electrodes. The dielectric spectra are described by the following equations:

$$\varepsilon' \approx \left(\frac{8q^2L^2}{\pi^4kT\varepsilon_0}\right) \times n_{ion}\left(\frac{1}{1+x^2} + \frac{1}{81+x^2} + \frac{1}{625+x^2}\right) + \varepsilon_s \tag{7}$$

$$\varepsilon'' \approx \left(\frac{8q^2L^2}{\pi^4kT\varepsilon_0}\right) \times n_{ion}\left(\frac{x}{1+x^2} + \frac{\frac{x}{9}}{81+x^2} + \frac{\frac{x}{25}}{625+x^2}\right)$$
 (8)

where  $x=\left(2L^2/\pi D\right)$ f,  $n_{ion}$  is the ionic concentration,q represents the electronic charge, D stands for the diffusion constant, L is the cell thickness, and T corresponds to the environmental temperature. Figure 4a,b present the experimental plots of  $\varepsilon'$  and  $\varepsilon''$  as a function of frequency for PPGDA monomers and PPGDA/E7 mixtures at room temperature (T = 24 °C). The fits in the range of 1–1000 Hz show good agreement with Equations (7) and (8), allowing us to extract different parameters such as  $n_{ion}$  and D. The spectra  $\varepsilon'$  and  $\varepsilon''$  were fitted together using Origin software. Tables 3 and 4 report the calculated values of D and  $n_{ion}$ .

**Table 3.** Static permittivity  $\varepsilon_s$ , diffusion constant D, and ionic concentration  $n_{ion}$  of the monomers at T = 24 °C (50  $\mu$ m thick samples).

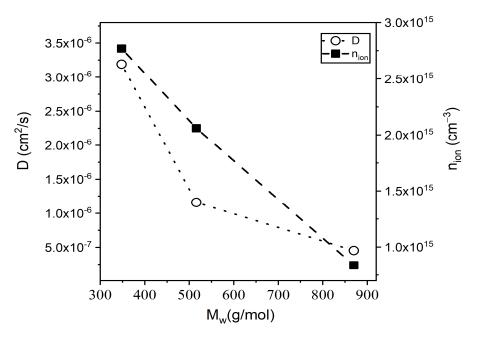
Monomer	$\epsilon_{\mathrm{s}}$	$n_{ion} \left( cm^{-3} \right) \left( \times 10^{15} \right)$	D(cm <sup>2</sup> /s) (×10 <sup>-6</sup> )
TPGDA	7.41	2.77	3.19
PPGDA540	6.47	2.06	1.16
PPGDA900	6.25	0.83	0.45

Crystals **2024**, 14, 286 11 of 15

<b>Table 4.</b> Static permittivity $\varepsilon_s$ , diffusion constant D, and ionic concentration $n_{ion}$ for monomer/E7
mixtures at T = 24 °C (50 $\mu$ m thick samples).

Monomer/E7 Mixture	$\epsilon_{ m s}$	$n_{ion}(cm-3) (\times 10^{13})$	D (cm <sup>2</sup> /s) (×10 <sup>-6</sup> )
70wt%TPGDA-30wt%E7	9.10	0.74	7.11
50wt%TPGDA-50wt%E7	10.30	1.99	7.80
30wt%TPGDA-70wt%E7	10.00	1.49	0.13
70wt%PPGDA540-30wt%E7	8.60	0.70	4.38
50wt%PPGDA540-50wt%E7	9.00	1.40	4.40
30wt%PPGDA540-70wt%E7	10.20	1.22	5.38
70wt%PPGDA900-30wt%E7	8.43	0.39	3.45
50wt%PPGDA900-50wt%E7	9.56	1.03	8.22
30wt%PPGDA900-70wt%E7	11.25	0.70	8.41

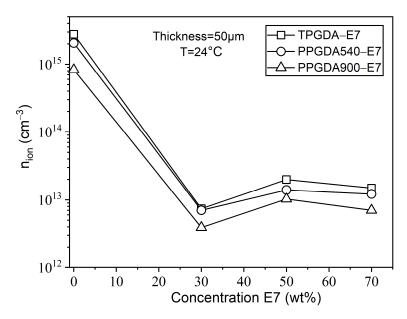
Figure 5 displays the relationship between the ionic concentration  $n_{ion}$  and the diffusion constant D, as a function of the molecular weight of the monomers at 24 °C. The plots show that as the molecular weight increases, both the ionic concentration and diffusion constant decrease. The decrease in the mobility of the monomer chain with higher molecular weight is likely responsible for the latter effect. The findings align with the thermal diffusion behavior observed in the binary aqueous solutions studied by Kishikawa et al. [44]. The relationship between D and the mobility (electrophoretic mobility)  $\eta$  is given by  $D = (kT/q)\eta$ . It can be inferred that the movement of the chain facilitates ionic migration towards the electrodes. The concentration of impurities in the monomers may be linked to the residues observed on the TGA analysis (inset of Figure 3a).



**Figure 5.** Evolution of the diffusion constant and ionic concentration as a function of molecular weight of the monomers at 24 °C.

Figure 6 displays the variation of ionic concentrations (n<sub>ion</sub>) in the monomer/LC mixtures as a function of the E7 concentration for the three PPGDA compounds. The decrease in the concentration of impurity observed when adding E7 to the three monomers can be attributed to the significantly lower ionic concentration of the pure LC compared to that of the monomers.

Crystals **2024**, 14, 286 12 of 15



**Figure 6.** Influence of the E7 concentration on the ionic concentration of the monomer/LC mixtures at 24 °C.

#### 4. Conclusions

This report investigates some of the thermophysical and dielectric responses of PPGDA monomers as a function of their molecular weight, as well as PPGDA mixed with nematic LCs E7. The molecular weights of the three monomers, TPGDA, PPGDA540, and PPGDA900, were determined using  $^1\text{H-NMR}$  analysis and found to be 350, 520, and 870 g.mol $^{-1}$ , respectively. These data represent average values, as the GPC analysis revealed some polydispersity effects on the molecular weights. The TGA analysis indicated that all monomers are thermally stable up to 100  $^{\circ}\text{C}$ . The residues at the end of each experiment (at 600  $^{\circ}\text{C}$ ) are likely to be of inorganic nature. DSC analysis of the monomers showed increasing Tg values as the molecular weight increased.

The static dielectric permittivity increases in the following order: PPGDA900, PPGDA540, and TPGDA. This can be explained by the higher dipole moment of TPGDA, which is caused by an enhanced volume density of carbonyl groups [7,8].

The electrical conductivity of the PPGDA and PPGDA/E7 blends was modeled using the Arrhenius and VTF models, both of which are applicable in describing the effects of conductivity in acrylic systems. The concentrations of ionic impurities and diffusion constants were estimated using the ionic hopping model. This approach shows that the ionic impurities and diffusion constants are inversely proportional to the molecular weight of PPGDA.

The investigation concludes that all monomer/LC blends contain ionic impurities, which can result in significant electrical conductivities. This may have a negative impact on the electro-optical properties of the resulting PDLC materials. The data presented in Figure 6 demonstrate a significant difference in ion concentrations between pure monomers and monomer/E7 blends. Consequently, the monomers have a much higher ionic content than E7. Therefore, it is essential to perform thorough purification procedures on the monomers to remove any ions before using them in PDLC materials.

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Crystals **2024**, 14, 286 13 of 15

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#### Nomenclature

$C_{m}$	Capacitance (F)
$E_a$	Activation energy (eV)
$E_{\mathbf{v}}$	Pseudo-activation energy (eV)
f	Frequency (Hz)
$G_{m}$	Conductance (S)
k	Boltzmann constant (k = $8.617 \times 10^{-5} \text{eV.K}^{-1}$ )
$M_{\rm w}$	Molecular weight $(g.mol^{-1})$
q	Electronic charge (q = $1.602 \times 10^{-19}$ C)
$\bar{T_{g}}$	Glass transition temperature (K)
$\varepsilon'(f)$	Relative permittivity spectrum
$\varepsilon''(f)$	Loss factor spectrum
$\varepsilon_0$	Permittivity of free space ( $\varepsilon_0 = 885 \times 10^{-14} \text{F.m}^{-1}$ )
σ	Conductivity (S.m <sup>-1</sup> )
$\sigma_1$	Pre-exponential factor $(S.K^{1/2}.m^{-1})$
TPGDA	Tri(propylene glycol) diacrylate (n = 3, CAS Number: 42978-66-5)
PPGDA540	Poly(propylene glycol) diacrylate (n = 7, CAS Number: 52496-08-9)
PPGDA900	Poly(propylene glycol) diacrylate (n = 12, CAS Number: 52496-08-9)

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